



U N I V E R S I T Y O F
SOUTH CAROLINA

LAMSS
Laboratory for Active Materials
and Smart Structures

A novel Helmholtz potential approach to predicting acoustic guided waves generated by fatigue crack

DOI: [10.18258/11075](https://doi.org/10.18258/11075)

Mohammad Faisal Haider
Mechanical Engineering Department
University of South Carolina

Acknowledgments

- SPARC, office of research, USC
- Crowdfunding (Experiment.com):
 - Farzana Yasmeen
 - Hanfei Mei
 - Amitav Tikadar
 - Jessica Livingston
 - Roshan Joseph
 - Lawrence Adrian
 - Matija Mazi
 - Md Rassel Raihan
 - Jon-Michael Adkins
 - Eric Frankforter
 - Amy Collette
 - Rodolfo Garcia-Contreras
 - LI Lingfang
 - WEEMING
 - Eric Damon Walters

Outline

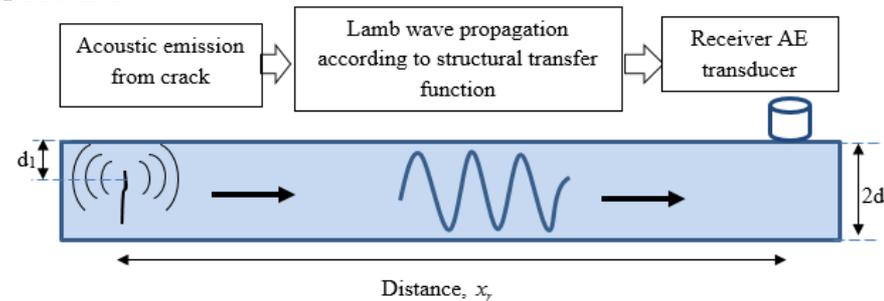
- Motivation
- Helmholtz potential approach to the analysis of acoustic emission (AE) guided wave
 - Theoretical formulation
 - Developing forward problem
- Developing inverse algorithm
- Experiments
- Summary, conclusions and future work

Objective

■ Acoustic Emission (AE) due to released energy

- Non-destructive testing (NDT)
- Locating and monitoring crack/ damage
- Crack characterization

■ Applications:



Nuclear spent fuel tank



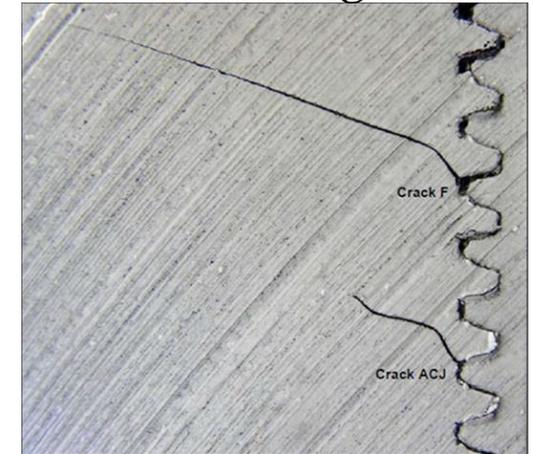
Ref: <http://www.sseb.org/downloads/Presentations/>

Aircraft fuselage



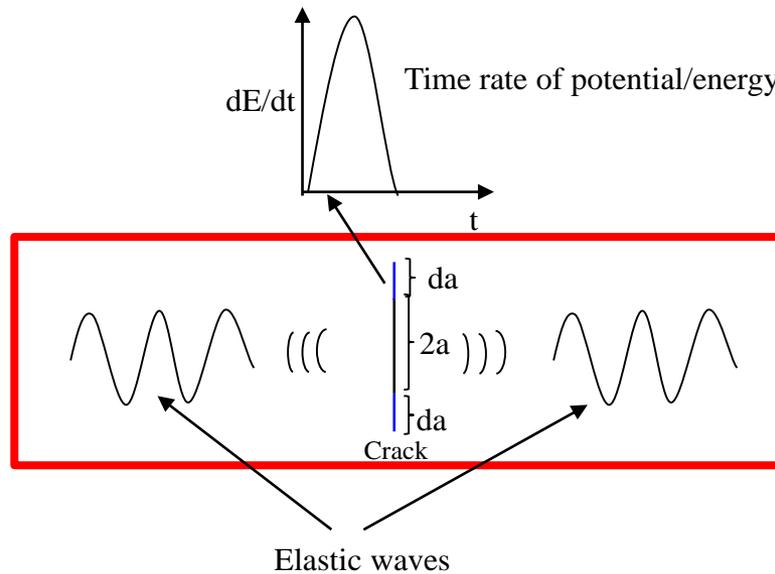
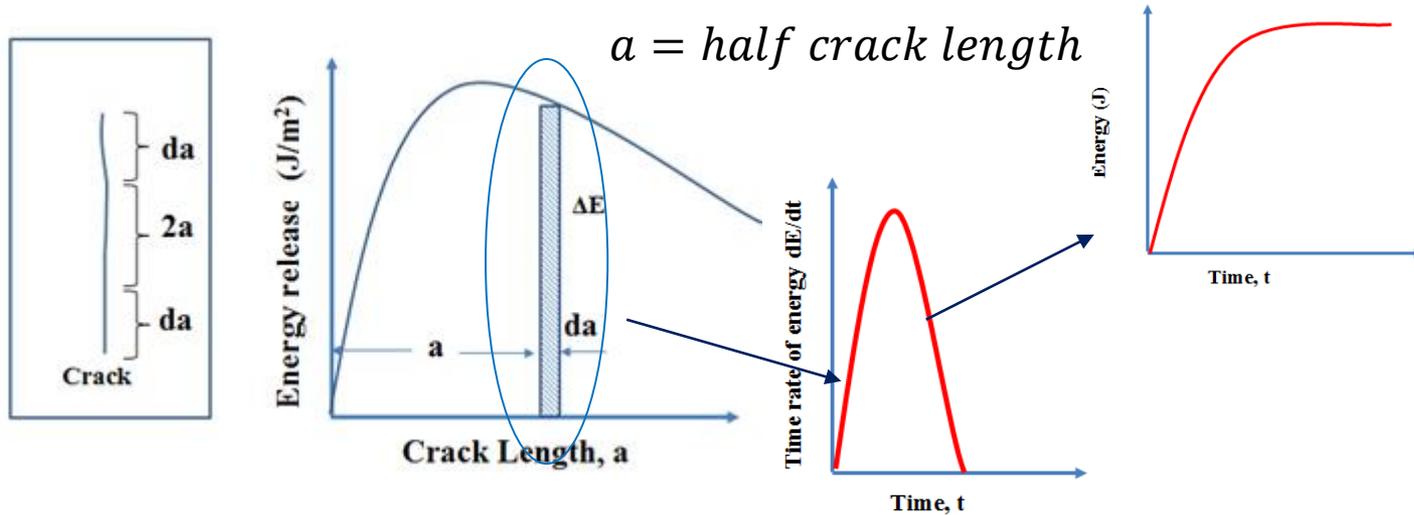
Ref: FSIMS Document

Gear teeth fatigue crack



Ref: <http://http://www.twi-global.com>

Overview of Potential Approach



Elastodynamic (Navier-Lame) Equations

■ Navier-Lame equations in vector form

$$(\lambda + \mu)\vec{\nabla}(\vec{\nabla} \cdot \vec{u}) + \mu\nabla^2\vec{u} = \rho\ddot{\vec{u}}$$

$$\vec{u} = u_x\hat{i} + u_y\hat{j} + u_z\hat{k}$$

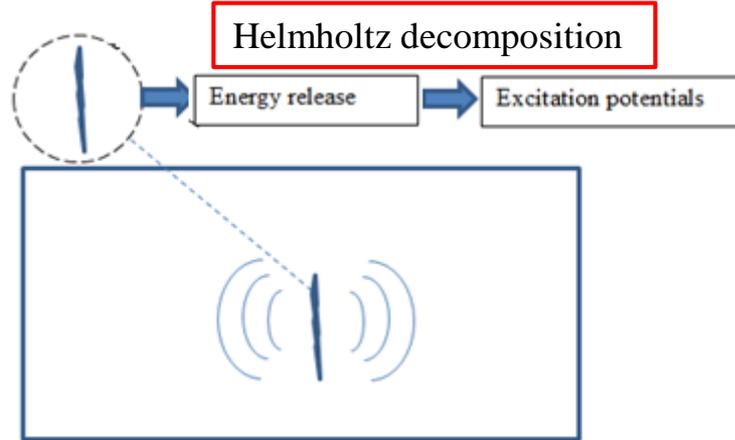
Lame constant λ, μ

Density, ρ

■ The basic concepts and equations of elastodynamics have been explained by

- Lamb (1917)
- Viktorov (1967)
- Graff (1991)
- Aki and Richards (2002)
- Achenbach (2003)
- Giurgiutiu (2014) etc.

Formation Pressure and Shear Excitation Potentials



Formation of scalar and vector excitation potentials

- If body force is present, then the Navier-Lame equation can be written as follow

$$(\lambda + \mu)\vec{\nabla}(\vec{\nabla} \cdot \vec{u}) + \mu\nabla^2\vec{u} + \rho\vec{f} = \rho\ddot{\vec{u}}$$

$$\vec{f} = f_x\hat{i} + f_y\hat{j} + f_z\hat{k}$$

$$\vec{u} = u_x\hat{i} + u_y\hat{j} + u_z\hat{k}$$

Lame constant λ, μ
Density, ρ

- Helmholtz decomposition states that any vector can be resolved into two potentials

$$\vec{\nabla} \cdot \vec{H} = 0 \quad \vec{H} = H_x\vec{i} + H_y\vec{j} + H_z\vec{k}$$

$$\vec{\nabla} \cdot \vec{B}^* = 0 \quad \vec{B}^* = B_x^*\vec{i} + B_y^*\vec{j} + B_z^*\vec{k}$$

$$\vec{u} = \text{grad } \Phi + \text{curl } \vec{H} = \vec{\nabla}\Phi + \vec{\nabla} \times \vec{H}$$

$$\vec{f} = \text{grad } A^* + \text{curl } \vec{B}^* = \vec{\nabla}A^* + \vec{\nabla} \times \vec{B}^*$$

Wave Equation for Potentials

- Inhomogeneous wave equation for potentials

$$\begin{aligned}
 c_P^2 \nabla^2 \Phi + A^* &= \ddot{\Phi} \\
 c_S^2 \nabla^2 \vec{H} + \vec{B}^* &= \ddot{\vec{H}} \longrightarrow \left\{ \begin{aligned} c_S^2 \nabla^2 H_x + B_x^* &= \ddot{H}_x \\ c_S^2 \nabla^2 H_y + B_y^* &= \ddot{H}_y \\ c_S^2 \nabla^2 H_z + B_z^* &= \ddot{H}_z \end{aligned} \right.
 \end{aligned}$$

- So for P+SV waves the relevant potentials are

$$\Phi, H_z, A^*, B_z^*$$

- Corresponding wave equations for potentials are

$$\begin{aligned}
 c_P^2 \nabla^2 \Phi + A^* &= \ddot{\Phi} \\
 c_S^2 \nabla^2 H_z + B_z^* &= \ddot{H}_z
 \end{aligned}$$

P+SV waves

Unknown displacement potentials (pressure and shear)

Excitation potentials (pressure and shear)

- Equations must be solved subject to zero-stress boundary conditions

Concentrated Potentials

$$\sigma_{xx}|_{y=\pm d} = 0, \sigma_{yy}|_{y=\pm d} = 0, \sigma_{xy}|_{y=\pm d} = 0$$

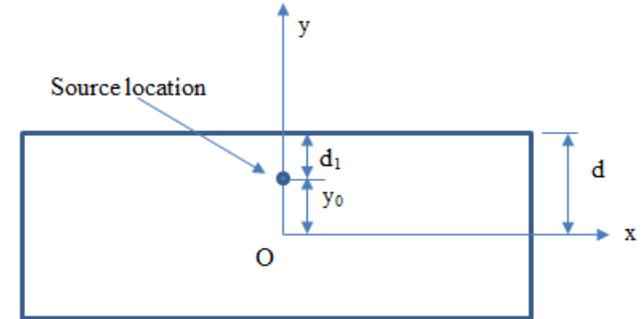
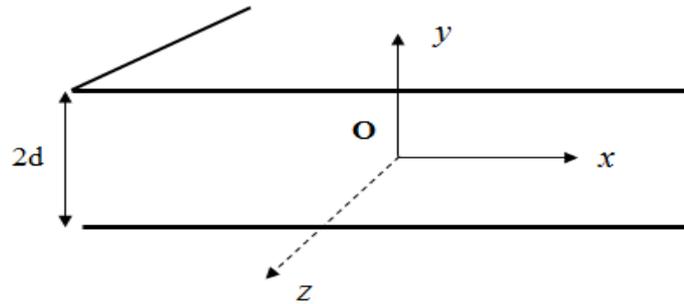


Plate of thickness $2d$ in which straight crested Lamb waves propagate in the x direction due to concentrated potential at $x = 0; y = 0$

- The potentials due to force can be written as follow

$$A = A\delta(x)\delta(y - y_0)e^{-i\omega t} \quad B_z = B_z\delta(x)\delta(y - y_0)e^{-i\omega t}$$

- Wave equations for potentials become

$$\frac{\partial^2 \Phi}{\partial x^2} + \frac{\partial^2 \Phi}{\partial y^2} + \frac{\omega^2}{c_p^2} \Phi = -A\delta(x)\delta(y - y_0)$$

$$\frac{\partial^2 H_z}{\partial x^2} + \frac{\partial^2 H_z}{\partial y^2} + \frac{\omega^2}{c_s^2} H_z = -B_z\delta(x)\delta(y - y_0)$$

$$c_p^2 \nabla^2 \Phi + A = \ddot{\Phi}$$

$$c_s^2 \nabla^2 \vec{H}_z + \vec{B}_z = \ddot{\vec{H}}_z$$

In-plane Strain Solution

Wave equations for potentials in a plate

Fourier transform in x direction

Thickness dependent differential equation

Total solution:

- (a) General solution for homogeneous equation
- (b) Particular solution due to source

Symmetric and antisymmetric Lamb wave solution

In-plane strain displacement is calculated

Lamb wave part

$$\varepsilon_x = i \left(\sum_{j=0}^{j_s} \left[P_S(\xi_j^S) \frac{N_S(\xi_j^S)}{D'_S(\xi_j^S)} \right] e^{i(\xi_j^S x - \omega t)} + \sum_{j=0}^{j_A} \left[P_A(\xi_j^A) \frac{N_A(\xi_j^A)}{D'_A(\xi_j^A)} \right] e^{i(\xi_j^A x - \omega t)} \right)$$

$$\begin{aligned} c_p^2 \nabla^2 \Phi + A &= \ddot{\Phi} \\ c_s^2 \nabla^2 \vec{H}_z + \vec{B}_z &= \ddot{\vec{H}}_z \end{aligned}$$

P+SV waves

$$\begin{aligned} \bar{\Phi}''(\xi, y) + \eta_p^2 \bar{\Phi}(\xi, y) &= -A \delta(y - y_0) \\ \bar{H}_z''(\xi, y) + \eta_s^2 \bar{H}_z(\xi, y) &= -B_z \delta(y - y_0) \end{aligned}$$

$$\begin{aligned} \bar{\Phi}(\xi, y) &= C_1 \sin \eta_p y + C_2 \cos \eta_p y - \frac{A}{2\eta_p} \sin \eta_p |y - y_0| \\ \bar{H}_z(\xi, y) &= i \left(D_1 \sin \eta_s y + D_2 \cos \eta_s y - \frac{B_z}{2\eta_s} \sin \eta_s |y - y_0| \right) \end{aligned}$$

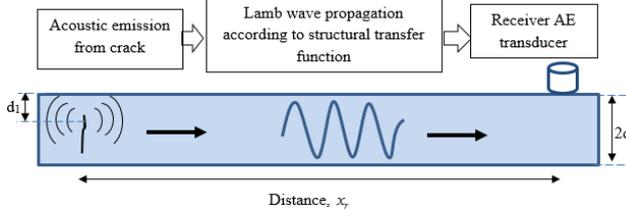
$$\begin{aligned} \begin{bmatrix} (\xi^2 - \eta_s^2) \cos \eta_p d & 2\xi \eta_s \cos \eta_s d \\ -2\xi \eta_p \sin \eta_p d & (\xi^2 - \eta_s^2) \sin \eta_s d \end{bmatrix} \begin{bmatrix} C_2 \\ D_1 \end{bmatrix} &= \begin{bmatrix} P_S \\ 0 \end{bmatrix} \\ \begin{bmatrix} (\xi^2 - \eta_s^2) \sin \eta_p d & -2\xi \eta_s \sin \eta_s d \\ 2\xi \eta_p \cos \eta_p d & (\xi^2 - \eta_s^2) \cos \eta_s d \end{bmatrix} \begin{bmatrix} C_1 \\ D_2 \end{bmatrix} &= \begin{bmatrix} 0 \\ P_A \end{bmatrix} \end{aligned}$$

$$P_S = \left((\xi^2 - \eta_s^2) \frac{A}{2\eta_p} \sin \eta_p d_1 + 2\xi \eta_s \frac{B_z}{2\eta_s} \cos \eta_s d_1 \right)$$

$$P_A = \left(2\xi \eta_p \frac{A}{2\eta_p} \cos \eta_p d_1 + (\xi^2 - \eta_s^2) \frac{B_z}{2\eta_s} \sin \eta_s d_1 \right)$$

Inverse Algorithm

Forward problem



Excitation Potentials:

Time rate of excitation potential $\dot{A}^*(t)$ and $\dot{B}_z^*(t)$
 Integrate: Time profile of released energy $A^*(t)$ and $B_z^*(t)$
 Calculate: Amplitude of the source
 $A(t) = \frac{h}{c_p} A^*(t)$; $B_z(t) = \frac{h}{c_s} B_z^*(t)$

Fourier transform of the signal

$$A(t) \rightarrow A(\omega); B_z(t) \rightarrow B_z(\omega)$$

Calculate source term $P_S(\xi^S)$ and $P_A(\xi^A)$

Calculate structural transfer function

$$\frac{N_S(\xi^S)}{D'_S(\xi^S)} e^{i(\xi^S x_r)} \text{ and } \frac{N_A(\xi^A)}{D'_A(\xi^A)} e^{i(\xi^A x_r)}$$

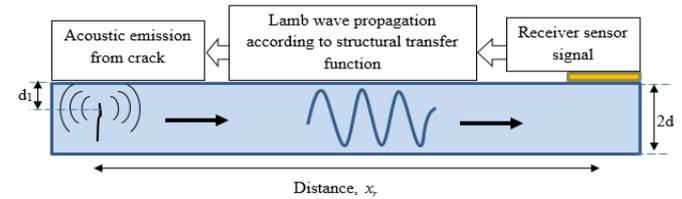
Product of structural transfer functions and source term

$$\varepsilon_x(\omega) = P_S(\xi^S, \omega) \frac{N_S(\xi^S)}{D'_S(\xi^S)} e^{i(\xi^S x_r)} + P_A(\xi^A, \omega) \frac{N_A(\xi^A)}{D'_A(\xi^A)} e^{i(\xi^A x_r)}$$

Perform inverse Fourier transform

$$\varepsilon_x(\omega) \rightarrow \varepsilon_x(t)$$

Inverse problem



Perform **Short time Fourier transforms (STFT)**

$$V_R(T) \rightarrow V_R(\omega)$$

Calculate structural transfer function

$$\frac{N_S(\xi^S)}{D'_S(\xi^S)} e^{i(\xi^S x_r)} \text{ and } \frac{N_A(\xi^A)}{D'_A(\xi^A)} e^{i(\xi^A x_r)}$$

Deconvolution of structural transfer functions and source term

$$P_S(\xi^S, \omega) = V_R^S(\omega) / \frac{N_S(\xi^S)}{D'_S(\xi^S)} e^{i(\xi^S x_r)}$$

$$P_A(\xi^A, \omega) = V_R^A(\omega) / \frac{N_A(\xi^A)}{D'_A(\xi^A)} e^{i(\xi^A x_r)}$$

Calculate source term $P_S(\xi^S)$ and $P_A(\xi^A)$

Calculate excitation source in frequency domain

$$A(\omega); B_z(\omega)$$

Inverse Fourier transform of the excitation signal

$$A(\omega) \rightarrow A(t); B_z(\omega) \rightarrow B_z(t)$$

Excitation Potentials:

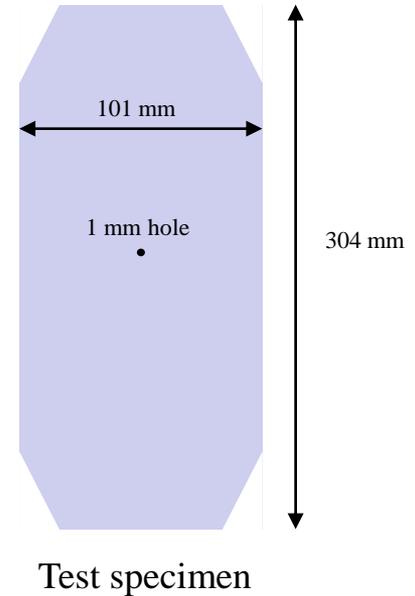
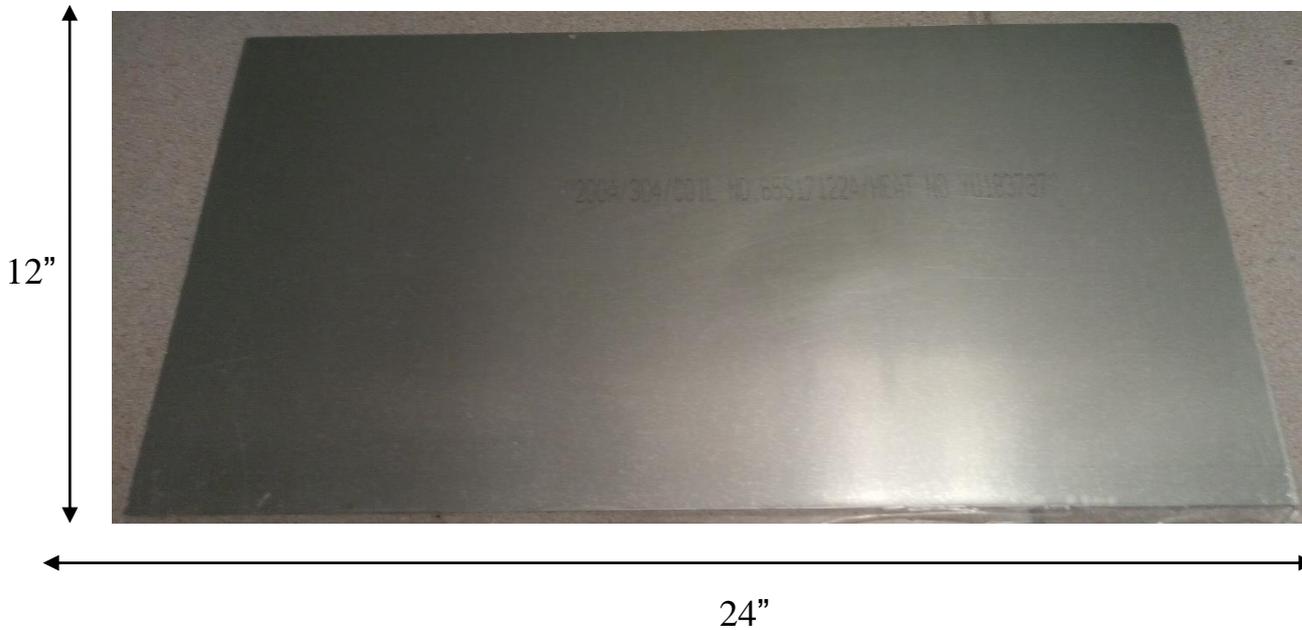
Calculate: Excitation Potentials source

$$A^*(t) = c_p^2 A(t); B_z^*(t) = c_s^2 B_z(t)$$

Experiments

Fabrication of Specimens

1 mm thick 304-stainless steel plate



Tensile Strength, Ultimate: 505 MPa
Tensile strength, Yield: 215 MPa

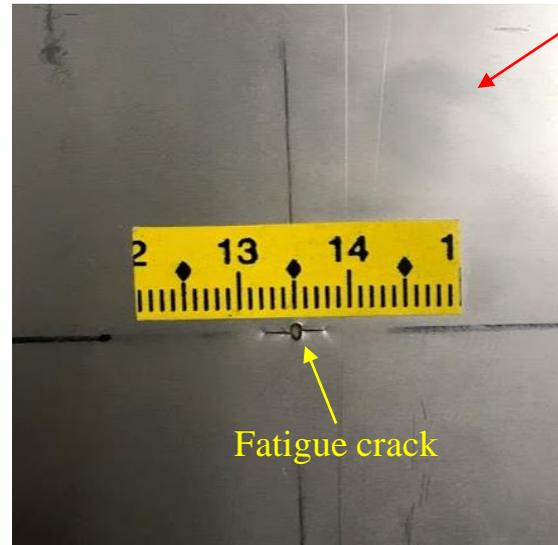
- 304 mm X 101 mm coupons are fabricated
- 1 mm hole is provided at the geometric center for crack initiation

Fatigue Crack Generation in Steel Specimen

- Fatigue loading is applied with a minimum load of 2.5 kN and maximum load of 25 kN at 4 Hz
- At this load range the crack was initiated after 20000 cycles
- After 35,000 fatigue cycles we observed 6 mm crack



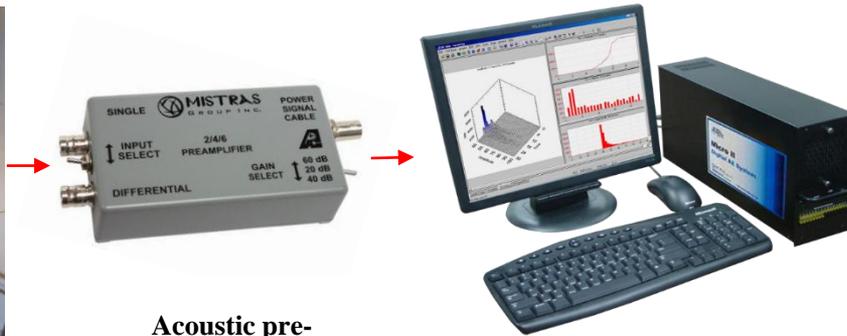
Fatigue test for crack generation



AE Measurement Experimental Setup

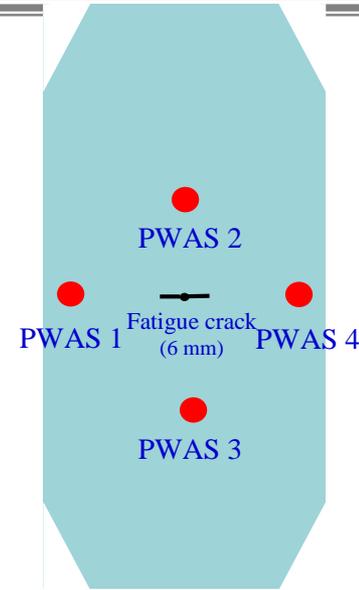


Test specimen mounted on
MTS machine



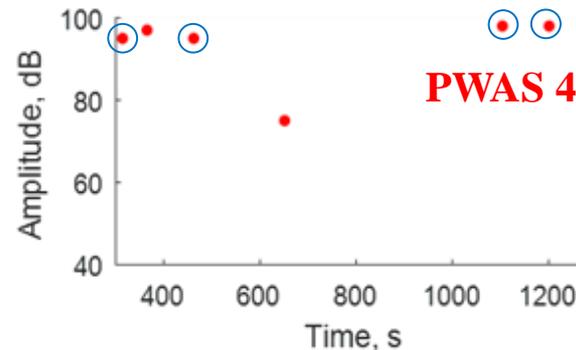
Acoustic pre-
amplifier

Mistras AE instrument



Test specimen

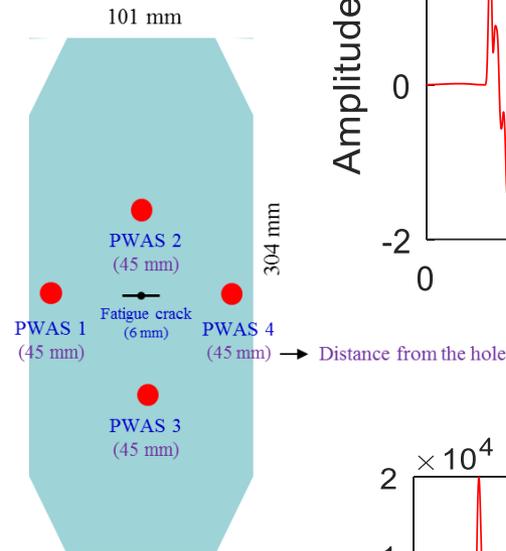
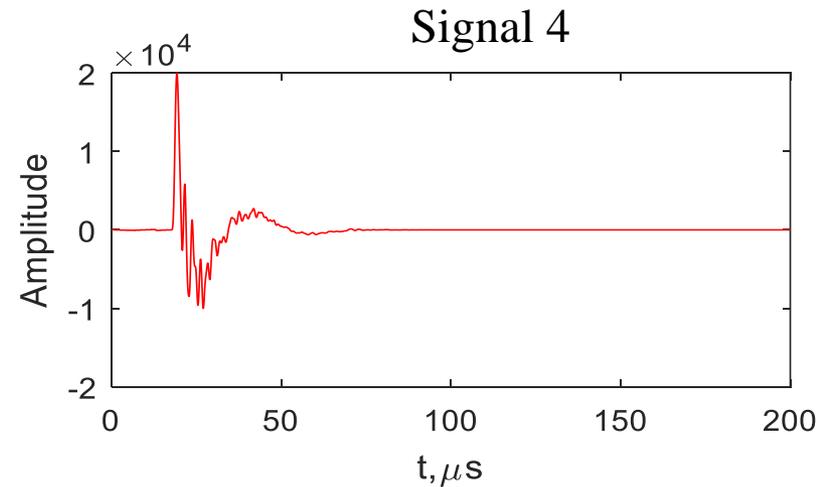
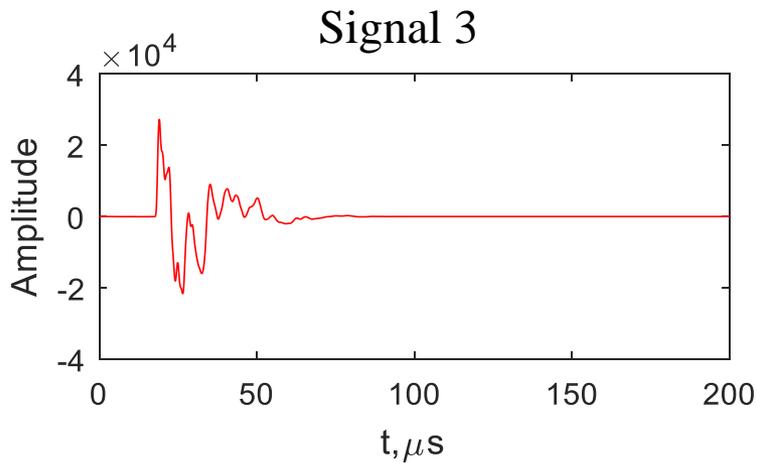
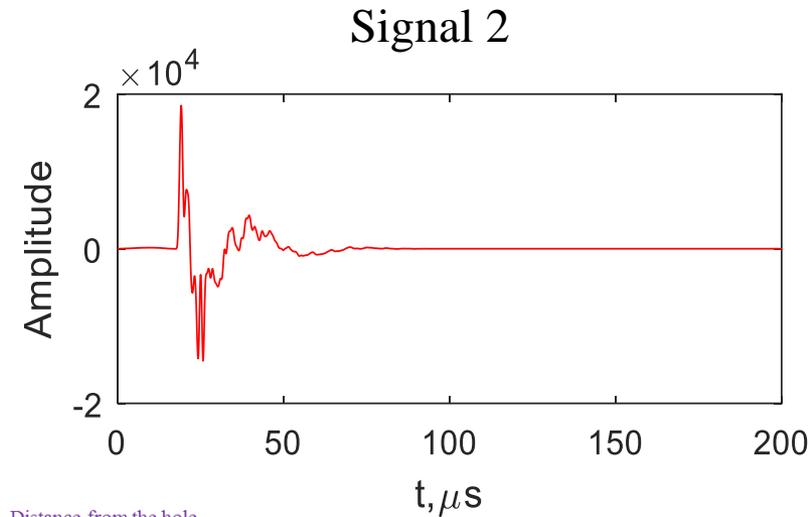
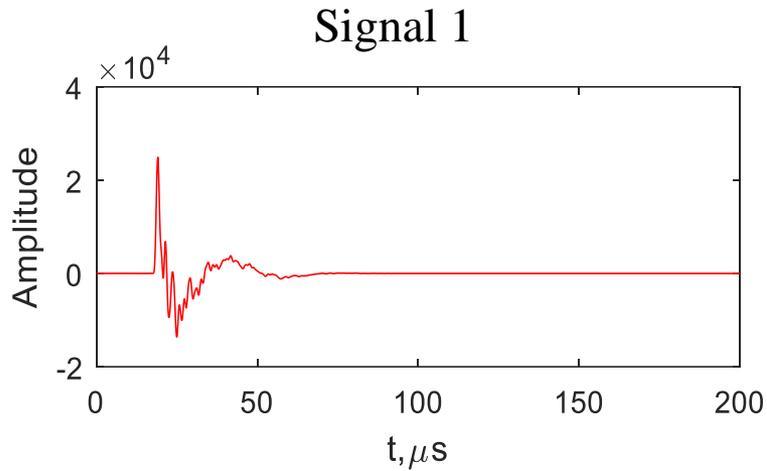
- Loading: 1.8 kN- 18 kN
- Frequency: 0.5 Hz
- Fatigue loading was applied
- AE signals were collected using Mistras AE system
- Approximate crack growth during fatigue experiment from 6 mm to 10 mm (5000 cycles)



AE hits at PWAS 4
over 600 cycles

Results

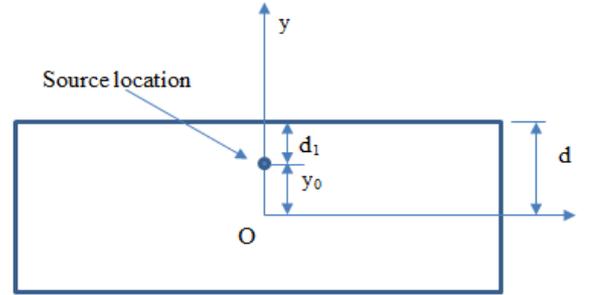
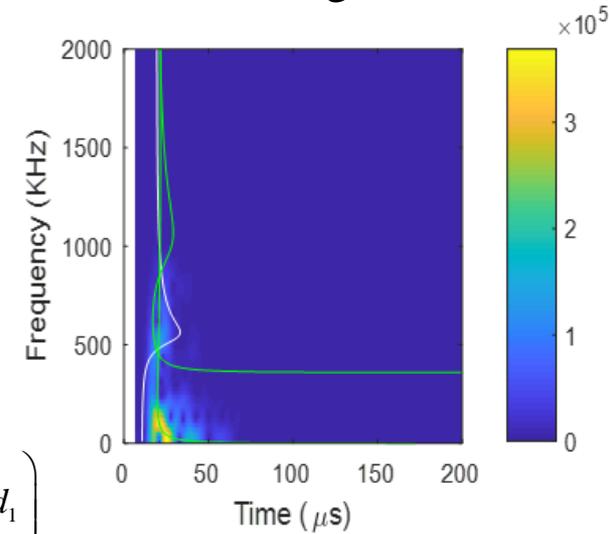
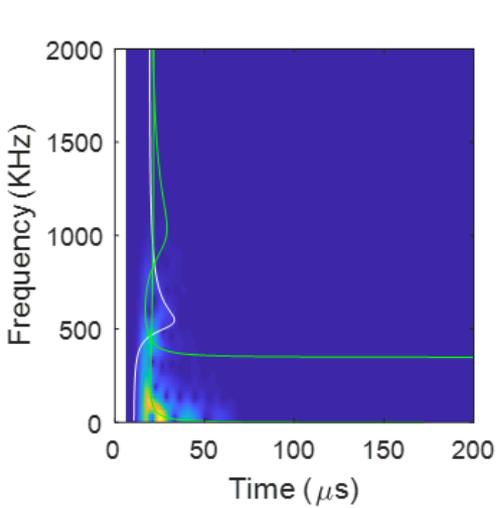
AE Signals Collected at PWAS 4



STFT of AE Signals

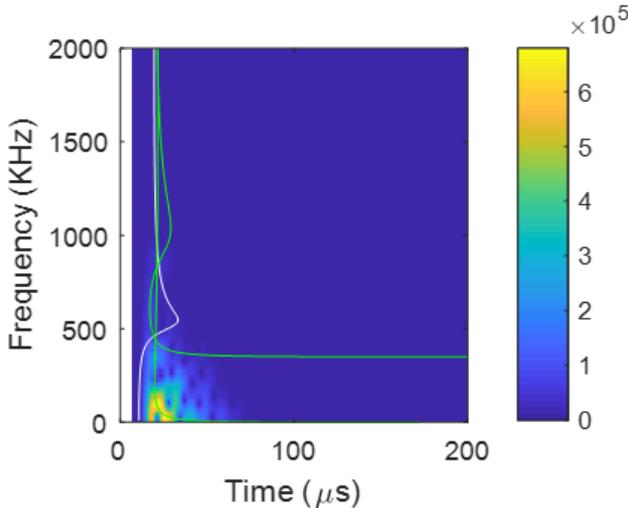
Signal 1

Signal 2

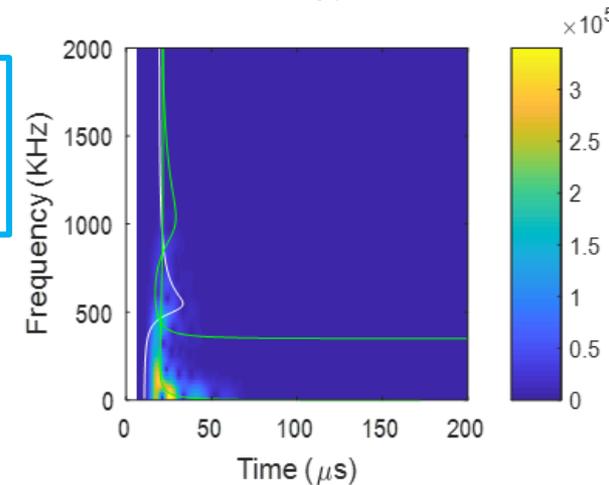


$$P_S = \left((\xi^2 - \eta_s^2) \frac{A}{2\eta_p} \sin \eta_p d_1 + 2\xi\eta_s \frac{B_z}{2\eta_s} \cos \eta_s d_1 \right)$$

Signal 4

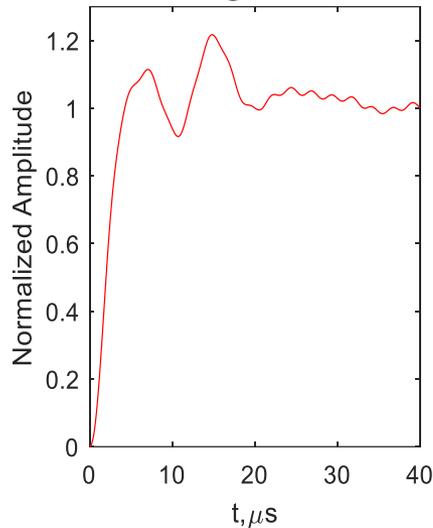


$$P_A = \left(2\xi\eta_p \frac{A}{2\eta_p} \cos \eta_p d_1 + (\xi^2 - \eta_s^2) \frac{B_z}{2\eta_s} \sin \eta_s d_1 \right)$$



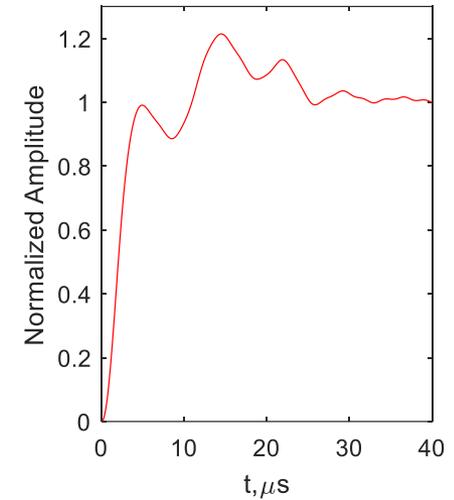
Pressure Excitation Potential Source of AE Signals

Signal 1

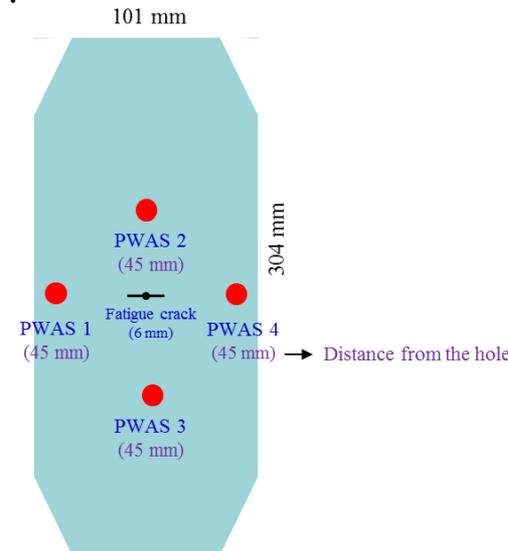


Rise time: $4.25 \mu\text{s}$

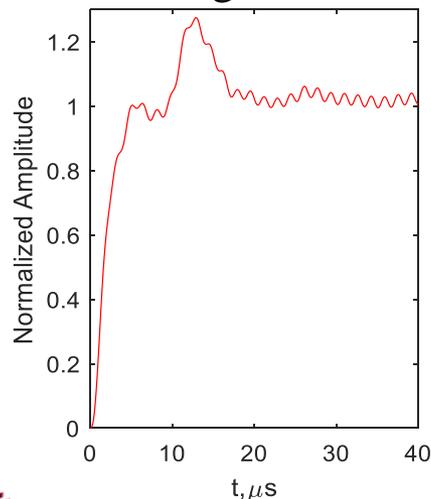
Signal 2



Rise time: $4.63 \mu\text{s}$



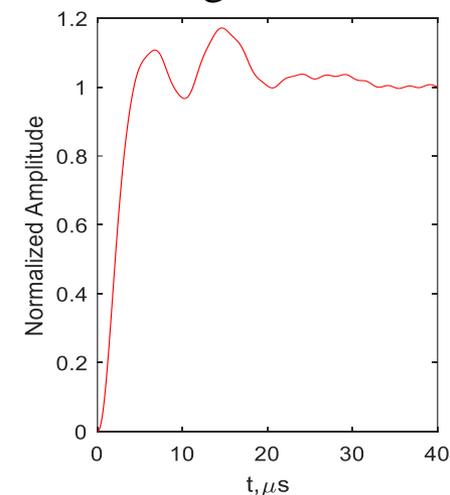
Signal 3



Rise time: $4.87 \mu\text{s}$

Rise time: $4.26 \mu\text{s}$

Signal 4



Summary and Conclusions

- The guided waves generated by an acoustic emission (AE) event were analyzed through a Helmholtz potential approach
- The inhomogeneous elastodynamic Navier-Lame equation was expressed as a system of wave equations in terms of
 - Unknown scalar and vector potentials
 - Known scalar and vector excitation potentials
- An experiment was designed and performed to extract AE signals from a fatigue crack growth
- An inverse algorithm was developed to characterize the AE source during crack propagation
- The source characterization can provide information
 - About excitation potential from the fatigue crack
 - About a qualitative as well as quantitative description of the crack propagation phenomenon

Thanks